



Wheeler–DeWitt Equation with Non-Minimal Coupling in a Closed FLRW Universe ($k = +1$)

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INTRODUCTION

Quantum cosmology describes the Universe through the Wheeler–DeWitt equation obtained through canonical quantization of the Hamiltonian constraint. In minisuperspace, a closed FLRW model reduces the problem to a few degrees of freedom, but non-minimal coupling ($\xi R\phi^2$) alters the kinetic structure and can change the hyperbolic character of the equation. Using Kiefer's canonical approach, we reduce the Wheeler–DeWitt equation to the closed FLRW case ($(k = +1)$, $(\Lambda = 0)$) and identify the conditions for hyperbolicity and the emergence of an intrinsic time variable.

MODEL

Consider the Einstein–Hilbert action with non-minimally coupled scalar field plus the Gibbons–Hawking–York edge term [1]

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2} (1 - \xi\phi^2) R - \frac{1}{2} g^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi - V(\phi) \right] + S_{\text{GHY}} \quad (1)$$

where the metric is considered as follows

$$ds^2 = -dt^2 + a^2(t) d\Omega_3^2, \quad (2)$$

with the spatial term such as

$$d\Omega_3^2 = d\chi^2 + \sin^2 \chi (d\theta^2 + \sin^2 \theta d\varphi^2), \quad (3)$$

$$0 \leq \chi \leq \pi. \quad (4)$$

After applying the ADM decomposition in the Hamiltonian formulation, a separation of kinetic terms must be carried out, for this purpose the redefinition of the field is used using the following ansatz

$$\chi = \sqrt{2\pi} \phi a^{6\xi} \quad (5)$$

REDUCED LAGRANGIAN

With the previous result, after applying the ansatz, the following Lagrangian is obtained that depends on the scale factor and the rescaling term of the scalar field $\mathcal{L}(a, \chi)$

$$\begin{aligned} L(a, \chi) = & -6\pi^2 a \dot{a}^2 + 6\pi^2 a + \frac{1}{2} a^{3-12\xi} \dot{\chi}^2 \\ & + 3\xi(1 - 6\xi) a^{1-12\xi} \dot{a}^2 \chi^2 \\ & - 3\xi a^{1-12\xi} \chi^2 - 2\pi^2 a^3 V\left(\frac{\chi}{\sqrt{2\pi} a^{6\xi}}\right). \end{aligned} \quad (6)$$

There are no mixed terms after the redefinition of the field. This is important for constructing the Lagrangian and subsequently finding solutions.

The effective coupling depends on χ and ξ

HAMILTONIAN

What is sought is for the Wheeler–DeWitt equation to have hyperbolic behavior; for this the Hamiltonian must meet the following condition

$$\mathcal{H} = 0 \quad , \quad \hat{p} \rightarrow -i\hbar\partial \quad (7)$$

$$[A(a, \chi) \partial_a^2 - B(a, \chi) \partial_\chi^2 + U(a, \chi)] \Psi(a, \chi) = 0 \quad (8)$$

The resulting action can be written as

$$S = \frac{1}{2} \int dt \left[-6a \dot{a}^2 \left(\frac{1}{48\pi G} + \xi \left(\xi - \frac{1}{6} \right) \chi^2 a^{-12\xi} \right) \right. \\ \left. + a^{3-12\xi} \dot{\chi}^2 + \frac{a}{8\pi G} - 6\xi \chi^2 a^{1-12\xi} - m^2 \chi^2 a^{3-12\xi} \right]. \quad (9)$$

and then calculating the conjugate moments in function of $\pi_a = \partial\mathcal{L}/\partial\dot{a}$ and $\pi_\chi = \partial\mathcal{L}/\partial\dot{\chi}$

WHEELER–DEWITT EQUATION

Substituting the above results into the Legendre transform to obtain the Hamiltonian with the restriction in terms of $(a, \chi; \pi_a, \pi_\chi)$

$$H = \frac{1}{2} \left[\frac{\pi_a^2}{6\xi(1 - 6\xi)\chi^2 a^{1-12\xi} - \frac{a}{8\pi G}} + a^{12\xi-3} \pi_\chi^2 \right] \\ - \frac{1}{2} \left[\frac{a}{8\pi G} - 6\xi \chi^2 a^{1-12\xi} - m^2 \chi^2 a^{3-12\xi} \right] = 0 \quad (10)$$

CONCLUSIONS

- The restricted Hamiltonian $H = 0$ leads to a Wheeler–DeWitt equation whose kinetic structure is governed by the “metric” of the minisuperspace (a, χ) . In particular, the sign of the gravitational component controls the type of equation (hyperbolic vs. elliptic).
- If $\xi = 1/12$, the type change is associated with a boundary defined by an algebraic condition at (a, χ) , which separates classically allowed and forbidden regions, and allows discussion of phenomena such as “phase space tunnelling”.

REFERENCES

[1] Valerio Faraoni. «Nonminimal Coupling». en. En: Fundamental Theories of Physics (2004). Ed. por Valerio Faraoni, págs. 143-195. DOI: [10.1007/978-1-4020-1989-0_7](https://doi.org/10.1007/978-1-4020-1989-0_7). (Visitado 01-12-2023).