



# Entropía De Entrelazamiento

Luis Carlos Desa Salas, 170735@unsaac.edu.pe

Escuela Profesional De Física, Universidad Nacional De San Antonio Abad Del cusco , Perú

## Introduction

In the current problem of the information paradox, it has led to generate new proposals for the entropy of black holes, in which the Hawking-Bekenstein entropy is incomplete, so a "generalized" entropy is proposed that has an extra term that considers quantum fields outside the event horizon [1]

$$S_{gen} = \frac{\mathcal{A}_{HS}}{4\hbar G_N} + S_{outside} \quad (1)$$

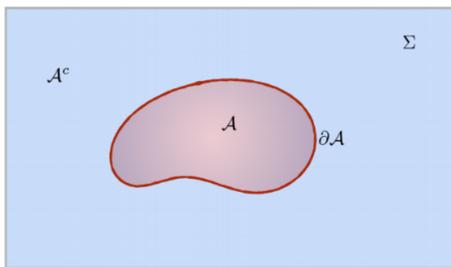
This proposal is based on the entropy proposed by Ryu-Takanayagi, this is the entropy of entanglement.

## Entanglement Entropy

The entangled state is defined as a whole, by the non-separability of the parts, therefore the pure quantum state belongs to the tensor product of the Hilbert space of the two subsystems, which are separated by the entanglement surface

$$|\psi\rangle \in H = H_A \otimes H_{A^c} \quad (2)$$

$$|\psi\rangle = \sum_{i,j} c_{ij} |i\rangle_A \otimes |j\rangle_{A^c} \quad (3)$$



In order to calculate the entanglement entropy we use the von Neumann equation, for which we need the reduced density matrix[3]

$$\rho_A = Tr_{A^c}(|\psi\rangle\langle\psi|) \quad (4)$$

$$S_A = -Tr_A(\rho_A \log \rho_A) \quad (5)$$

If the reduced density matrix is diagonalizable

$$S_A = -\sum(\lambda_i \log \lambda_i) \quad (6)$$

## Continuum QFTs

To calculate the entanglement entropy of continuous fields, we must make a temporary slice ( $t = 0$ ), that for a d-dimensional QFT, ( $\Sigma_{d-1}$ ), the reduced density matrix will be a path integral with the boundary conditions of the past and future causal of our temporal slice

$$(\rho_A)_{-+} = \int [D\Phi] e^{-S_{QFT}[\Phi]} \delta_E(\Phi_{\mp A}) \quad (7)$$

$$\delta_E(\Phi_{\mp A}) = \delta(\Phi_A(t=0^- - \Phi_-)) \delta(\Phi_A(t=0^+ - \Phi_+)) \quad (8)$$

## Entanglement Entropy of $CFT_3$

With the considerations mentioned above, we can obtain the following result for a d-dimensional QFT, where the first term is the area law, it is required that ( $d \geq 3$ ) [2]

$$S_{EE}^{QFT} = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_1 \frac{R^{d-4}}{\delta^{d-4}} + \dots + a \log \frac{R}{\delta} - F \quad (9)$$

For a circumference of radius  $R$ , that is ( $d = 3$ ), we obtain the following result, (it can also be for non-smooth surfaces)

$$S_{EE} = c_0 \frac{R}{\delta} - F \quad (10)$$

Where ( $\delta$ ) is a k-top, a scale constant.

## Holographic Entanglement Entropy

Inspired by the area law, the following equation is proposed to calculate the holographic entanglement entropy

$$S_{EE}(A) = \frac{\min Area(\gamma_A)}{4G_N} \quad (11)$$

Where ( $\gamma_A$ ) is a surface that covers the circumference and extends into space-time  $AdS_4$ , the minimum value of this surface is used, then ( $Area(\gamma_A)$ ) is an extremal. For the calculation we will use the Poincare Patch metric

$$ds_{PP}^2 = \frac{L^2}{z^2} [-dt^2 + dz^2 + dr^2 + r^2 d\theta^2] \quad (12)$$

With  $z = f(r)$  and  $\dot{f}(r) = df(r)/dr$ , how ( $t = 0$ )  $\Rightarrow (dt = 0)$

$$ds_{\gamma}^2 = \frac{L^2}{f(r)^2} [(\dot{f}(r)^2 + 1)dr^2 + r^2 d\theta^2] \quad (13)$$

With the determinant of its metric ( $\det(h) = \sqrt{h}$ )

$$\sqrt{h} = \frac{L^2}{f(r)^2} r \sqrt{1 + \dot{f}(r)^2} \quad (14)$$

The entanglement entropy

$$S_{EE}(A) = \frac{1}{4G} \min \int \sqrt{h} dr d\theta \quad (15)$$

$$S_{EE}(A) = \frac{1}{4G} \min \int \frac{L^2}{f(r)^2} r \sqrt{1 + \dot{f}(r)^2} dr d\theta \quad (16)$$

Applying the Euler-Lagrange equations, inside the integral, the function  $f(r)$  that satisfies it is ( $z = f(r) = \sqrt{R^2 - r^2}$ ), replacing

$$S_{EE}(A) = \frac{\pi L^2}{2G} \int_0^R \frac{r}{(R^2 - r^2)^{3/2}} dr \quad (17)$$

Solving the integral, we obtain the following result, which is the same as the one obtained previously, with a quantum calculation (10).

$$S_{EE}(A) = \frac{\pi L^2 R}{2G \delta} - \frac{\pi L^2}{2G} \quad (18)$$

$$S_{EE}(A) = c_0 \frac{R}{\delta} - F \quad (19)$$

## Conclusions

- Thus we can conclude that the holographic calculation with the equation (11) for the entanglement entropy is much more efficient than from a quantum theory  $CFT$ .

## References

- [1] M. S. T. Almheiri, Hartman. The entropy of hawking radiation, 2020. URL [arXiv:2006.06872v1](https://arxiv.org/abs/2006.06872v1) [hep-th].
- [2] C. Calabrese. Entanglement entropy and conformal field theory, 2009. URL [arXiv:0905.4013v2](https://arxiv.org/abs/0905.4013v2) [cond-mat.stat-mech].
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